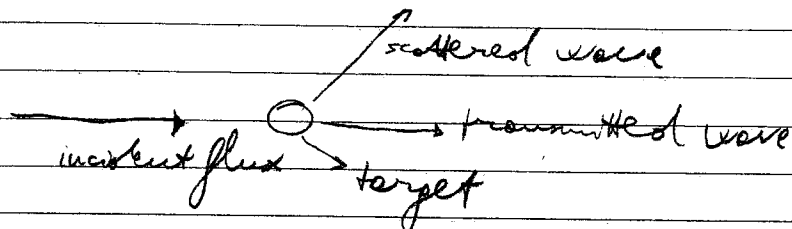


Lecture 20



Elements of scattering theory

Scattering cross section:



Assume the Hamiltonian is

$$H = H_0 + V(\vec{r})$$

$$H_0 \equiv \frac{-\hbar^2 \nabla^2}{2m}$$

$V(\vec{r}) \rightarrow 0$ as $r \rightarrow \infty$ more rapidly than $\frac{1}{r}$

Before the collision at $t \rightarrow -\infty$ a particle can be described by the wave function (along the z direction)

$$\psi_k = e^{ikz}$$

Define two stationary waves

$$\psi_k^{(+)} = e^{ikz} + \text{outgoing wave}$$

$$\psi_k^{(-)} = e^{-ikz} + \text{incoming wave}$$

If E is the energy of the particle

$$k = \sqrt{2mE}$$

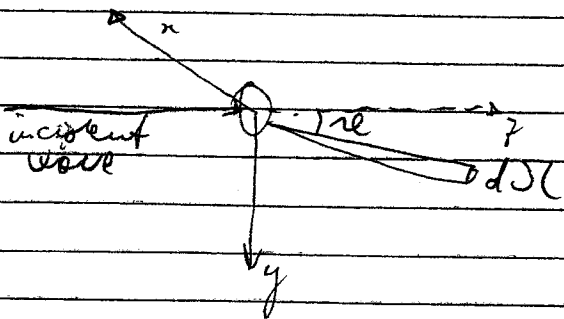
$\psi_k^{(+)}$ is solution to

$$\# \psi_k^{(+)} = E \psi_k^{(+)}$$

with boundary condition at $r \rightarrow \infty$

$$\psi_k^{(+)} \rightarrow e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$

$f(\theta)$ is called scattering amplitude



The number of particles scattered into the vector solid angle $d\Omega$ is given by the radial component of the scattered current density 3

$$\begin{aligned} \vec{J}_{sc,r} &= \frac{\hbar}{2mi} \left(\left(\psi_k^* \right) \frac{\partial \psi_k^{(sc)}}{\partial r} - \psi_k^{(sc)} \frac{\partial \psi_k^*}{\partial r} \right) = \\ &= \frac{\hbar k}{mr^2} |f(\theta)|^2 \end{aligned}$$

Number of particles passing through the surface $d\vec{S}$ (defined by the angle $d\Omega$) per unit time

$$dN = \vec{J}_{sc} \cdot d\vec{S} = r^2 \vec{J}_{sc} \cdot d\Omega = r^2 \vec{J}_{sc,r} d\Omega$$

dN has to be proportional to \vec{J}_{inc}

$$dN \propto \vec{J}_{inc} \Rightarrow dN = d\sigma \vec{J}_{inc}$$

$d\sigma$ is the differential scattering cross section

The total scattering cross section

$$\sigma = \int d\sigma = \int_{4\pi} \left(\frac{d\sigma}{d\Omega} \right) d\Omega$$

$$r^2 J_{sc,r} d\Omega = J_{inc} d\Omega$$

$$J_{inc} = \frac{I_0 k}{4\pi} \Rightarrow r^2 J_{sc,r} d\Omega = \frac{I_0 k}{4\pi} d\Omega$$

⇓

$$d\Omega = |f(\alpha)|^2 d\alpha$$

Total cross section

$$\sigma = \int d\sigma = 2\pi \int_0^\pi |f(\alpha)|^2 \sin \alpha d\alpha$$