

lecture 14

Solution of the time-dependent Schrödinger equation

(1) $i\hbar \frac{d\psi(t)}{dt} = \hat{H} \psi(t)$ [let us assume to begin with that V in \hat{H} does not depend on time. We'll generalise later]

with initial condition $\psi(0) = \psi_0$

linearity of (1) $\Rightarrow \psi(t) = \hat{U}(t) \psi_0$

$\hat{U}(t)$ is a linear operator: operator of time evolution

$\langle \hat{U}(t) \psi_0 | \hat{U}(t) \psi_0 \rangle = \langle \psi_0 | \psi_0 \rangle = 1$ [isometric and continuous operator]

$\hat{U}(0) = 1$; $\hat{U}(t_1) \hat{U}(t_2) = \hat{U}(t_1 + t_2)$ [only if \hat{H} does not dep. on time]

$\psi(t) = \hat{U}(t - t_0) \psi_0$ [this shows that if $V(x)$ does not depend on time, t_0 is not privileged]

$$\psi(x, t) = \sum_i c_i \psi_i(x, t) + \int d\mathcal{B} c(\mathcal{B}) \psi_{\mathcal{B}}(x, t) =$$

$$= \sum_i c_i u_i(x) e^{-\frac{i\mathcal{E}_i t}{\hbar}} + \int d\mathcal{B} c(\mathcal{B}) u_{\mathcal{B}}(x) e^{-\frac{i\mathcal{E}_{\mathcal{B}} t}{\hbar}}$$

From (1)

[2]

$$i\hbar \frac{d}{dt} \hat{U}(t) \psi_0 = \hat{H} \hat{U}(t) \psi_0$$

Since ψ_0 is arbitrary

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$$(2) \quad i\hbar \frac{d}{dt} \hat{U}(t) = \hat{H} \hat{U}(t) \quad \text{to be solved with initial condition}$$
$$\hat{U}(0) = 1$$

Formal solution

$$\hat{U}(t) = e^{-\frac{i\hat{H}t}{\hbar}} \quad [\hat{H} \text{ indep. of time}]$$

(2) is useful when \hat{H} depends on time

Rewrite (2) in integral form with $\hat{U}(0) = 1$

$$(3) \quad \hat{U}(t) = 1 - \frac{i}{\hbar} \int_0^t dt_1 \hat{H}(t_1) \hat{U}(t_1)$$

Iterate the solution (3) and hence the

Assume $\hat{U}(t_1)$ assumes the unperturbed value 3

$$\hat{U}(t_0) = 1$$

$$\hat{U}(t) \approx 1 - \frac{i}{\hbar} \int_0^t dt_1 \hat{H}(t_1)$$

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$$U(t) \approx 1 - \frac{i}{\hbar} \int_0^t dt_1 \hat{H}(t_1) \left[1 - \frac{i}{\hbar} \int_0^{t_1} dt_2 \hat{H}(t_2) \right]$$

$$= 1 - \frac{i}{\hbar} \int_0^t dt_1 \hat{H}(t_1) + \left(-\frac{i}{\hbar} \right)^2 \int_0^t dt_1 \int_0^{t_1} dt_2 \hat{H}(t_1) \hat{H}(t_2)$$

If I iterate this equation, i.e. put this solution back into (3) I obtain

$$\hat{U}(t) = \sum_{r=0}^{\infty} \left(\frac{-i}{\hbar} \right)^r \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{r-1}} dt_r \hat{H}(t_1) \hat{H}(t_2) \dots \hat{H}(t_r)$$

Define $T[\hat{H}(t_1) \hat{H}(t_2) \dots \hat{H}(t_r)]$ chronological ordered product of \hat{H} s

$$T[\hat{H}(t_1) \hat{H}(t_2)] = \begin{cases} \hat{H}(t_1) \hat{H}(t_2) & \text{if } t_1 \geq t_2 \\ \hat{H}(t_2) \hat{H}(t_1) & \text{if } t < t_2 \end{cases}$$

invariant with respect of permutation of times

$$\hat{U}(t) = \sum_{r=0}^{\infty} \frac{1}{r!} \left(-\frac{i}{\hbar} \right)^r \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{r-1}} dt_r T \left[\hat{H}(t_1) \hat{H}(t_2) \dots \hat{H}(t_r) \right] \quad \boxed{4}$$

$$\hat{U}(t) = T \left\{ e \left[-\frac{i}{\hbar} \int_0^t dt' \hat{H}(t') \right] \right\}$$

Dyson expansion. If \hat{H} does not depend on t

$$\hat{U}(t) = e^{-\frac{i\hat{H}t}{\hbar}}$$

In general $\hat{U}(t, t_0) = \hat{U}(t) \hat{U}^{-1}(t_0) = T \left\{ e \left[-\frac{i}{\hbar} \int_{t_0}^t dt' \hat{H}(t') \right] \right\}$
 $\times U(t-t_0)$