

lecture 5



Identical fermions

Consider an  $N$ -fermion system  
due to the Pauli exclusion principle

the wave function of the  $N$ -particle system with  <sup>$N$</sup>  single particle state  $|\psi_1\rangle, |\psi_2\rangle, |\psi_3\rangle, \dots$  can be written as a  $N \times N$  Slater determinant

$$S = \frac{1}{\sqrt{N!}} \begin{vmatrix} |\psi_1\rangle & |\psi_2\rangle & \dots & |\psi_N\rangle \\ |\psi_1\rangle & |\psi_2\rangle & \dots & |\psi_N\rangle \\ \vdots & \vdots & \ddots & \vdots \\ |\psi_N\rangle & |\psi_N\rangle & \dots & |\psi_N\rangle \end{vmatrix}$$

So to each set  $|\psi_1\rangle, |\psi_2\rangle, \dots, |\psi_N\rangle$  of  $N$  different states chosen from among the individual states there corresponds one and only one, antisymmetrized state, represented by  $S$

These states constitutes an orthonormal basis of a subspace  $\mathcal{S}^{(N)}$  of the space  $\mathcal{S}$  of all dynamical states of the whole system of  $N$  generic particles. [2]

There are  $\binom{d_n}{N_n}$  ways to choose the  $N_n$  occupied states in the single-particle energy  $E_n$

$$Q(N_1, N_2, N_3, \dots) = \prod_{n=1}^{\infty} \frac{d_n!}{N_n! (d_n - N_n)!}$$

At thermal equilibrium <sup>and weak interactions</sup> every state <sup>many-particle</sup> with a given total energy  $E$  and  $N$  particles is equally probable.

The most probable configuration  $(N_1, N_2, \dots)$  is the one can be achieved with largest number of ways subject to the constraints:

$$\sum_{n=1}^{\infty} N_n = N$$

$$\sum_{n=1}^{\infty} N_n E_n = E$$

[3]

Solve the maximization procedure:

$$N_n = \frac{d_n}{e^{(\alpha + \beta E_n)} + 1}$$

$\alpha, \beta$  are Lagrange multipliers that play the role of chemical potential

$$\mu(T) = -\alpha k_B T$$

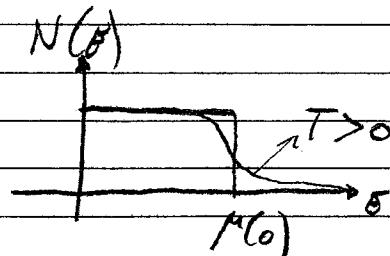
and  $\beta = \frac{1}{k_B T}$

So

$$N_n = \frac{d_n}{e^{(E_n - \mu)/k_B T} + 1}$$

Fermi-Dirac distribution

For  $d_n = 1$   $E_n \rightarrow E$   $N(E) \rightarrow$



## Identical bosons

4

Assume a gas of weakly interacting bosons.

Symmetrization of wavefunctions require one  $N$ -particle state where a specific set of 1-particle states are occupied.

No restriction on the number of particles with same 1-particle state.

Solution

$$Q(N_1, N_2, \dots) = \prod_{u=1}^{\infty} \frac{(N_u + d_u - 1)!}{N_u! (d_u - 1)!}$$

Performing the constrained maximization:

$$N_u = \frac{d_u}{e^{(\epsilon_u - \mu)/k_B T}} \quad \text{Bose-Einstein}$$

For classical particles

$$N_u = \frac{d_u}{e^{(\epsilon_u - \mu)/k_B T}} \quad \text{Maxwell-Boltzmann}$$